

## RAYLEIGH NUMBER DEPENDENCE OF THE ARCHIMEDES NUMBER DEPENDENT LARGE-SCALE FLOW STRUCTURE FORMATION IN MIXED CONVECTION

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**Abstract** We report on the experimental investigations of large-scale flow structure formation in mixed convection. We characterize the flow field by measuring the velocity fields within a rectangular model room using 2D2C PIV. The control parameters are the Reynolds number  $Re$ , the Rayleigh number  $Ra$  and the Prandtl number  $Pr$ . All parameters are linked through the Archimedes number  $Ar$ . In  $6.4 \times 10^{-2} \leq Ar \leq 1.39 \times 10^1$ ,  $4.2 \times 10^3 \leq Re \leq 6.35 \times 10^4$  and  $Ra = 3.1 \times 10^7$ ,  $Ra = 1.8 \times 10^8$  and  $Pr = 0.713$  we found flow 3 different flow structures. While keeping  $Ra$  and  $Pr$  constant and varying  $Ar$  through  $Re$  variations, we found an  $Ar$  dependence of the large-scale flow structure formation within  $6.4 \times 10^{-2} \leq Ar \leq 1.39 \times 10^1$ . Furthermore, we found a  $Ra$  dependence of the structure formation, which shifts the transition points between the structures to higher Archimedes numbers and reduces the mean velocities within the investigated domain.

### INTRODUCTION

Mixed convection is defined as the temporal and spatial superposition of natural and forced convection and is driven by inlet velocity  $u_{in}$  and a temperature gradient  $\Delta T$  within a gravity field. In most applications, e.g. indoor airflows, this flow is turbulent and its temperature and velocity field is dominated by the formation of large-scale flow structures, as comprehensively studied in Rayleigh-Benard-Convection (RBC) (review of RBC see e.g. [1]). These structures verifiable have a high impact on matter and heat transport properties in the considered domain [2, 3]. The control parameters  $u_{in}$  and  $\Delta T$  can be expressed by the Reynolds number  $Re$ , the Rayleigh number  $Ra$  the Prandtl number  $Pr$  (see [4] for definitions). The relation between them is defined as Archimedes number  $Ar$ , which seems to be the main control parameter for the interaction of forced and natural convection ( $Ar < 0.1$  ... forced convection;  $Ar \approx 1$  ... mixed convection;  $Ar \gg 1$  natural convection). Beyond the found large-scale flow structures, we aim to show that the Archimedes dependent structure formation also depends on the Rayleigh number  $Ra$ . This might have a strong impact on the predictability of the large-scale flow structures.

### EXPERIMENTAL SETUP

Figure 1a shows the model room in Cartesian coordinates  $x, y, z$ , which are subsequently normalized to  $L, H, D$ , respectively. The room has a rectangular shape with aspect ratios  $\Gamma_{xy} = L/H = 1.33$  and  $\Gamma_{xz} = D/H = 1.66$  and a height of  $H = 0.3$ m. The cross-section shown in Fig. 1a is constant along the depth of the room. The global temperature difference  $\Delta T$  is defined as difference of the mean temperature of the four electrical heated obstacles at the bottom and the mean temperature of the air inlets at the top of the side walls. The air inlets and outlets at top and bottom of the side walls are generating the inlet velocity  $u_{in}$ . For more details on room and boundary conditions, please refer to [4, 5].

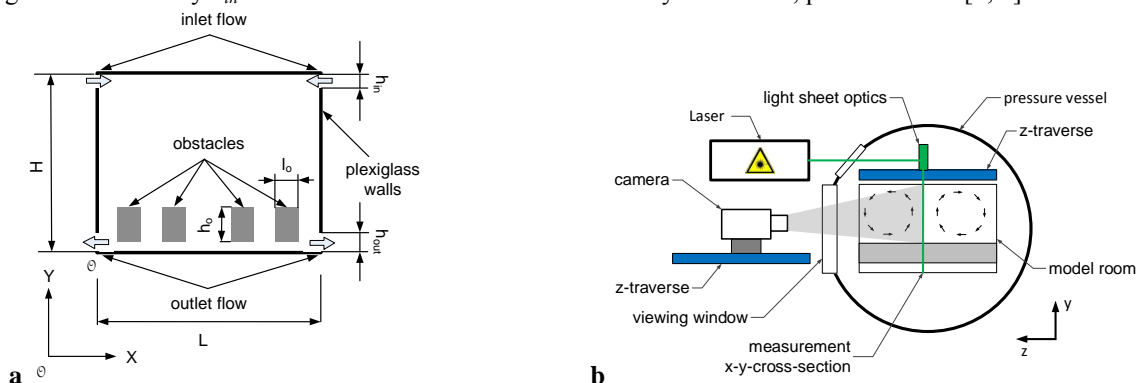


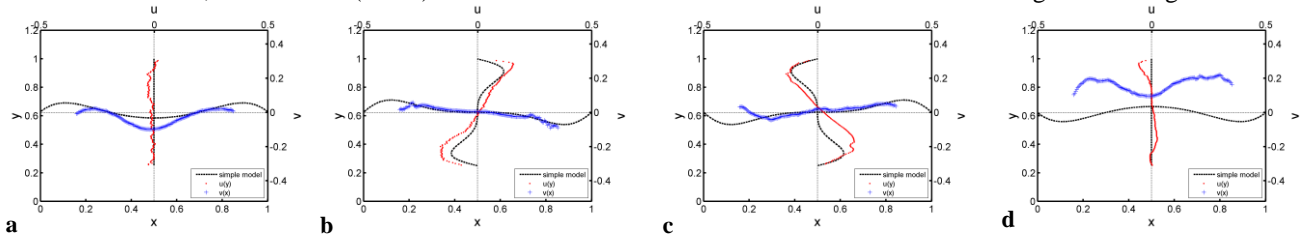
Fig. 1 a Schematic of the investigate geometry (model room) [5]; b schematic of the experimental setup

Two Rayleigh numbers have been investigated  $Ra = 3.1 \times 10^7$  and  $Ra = 1.8 \times 10^8$  at constant Prandtl number  $Pr = 0.713$  and  $\Delta T = 10$ K. The Rayleigh numbers have been achieved by using compressed air at pressures  $p = 1.0 \times 10^5$ Pa and  $p = 2.51 \times 10^5$ Pa. Thus, the experiments are carried out in a pressure vessel, as shown in Fig. 1b. In order to achieve an Archimedes number variation in the range of  $6.4 \times 10^{-2} \leq Ar \leq 1.39 \times 10^1$ , the Reynolds numbers are varied between  $6.82 \times 10^3 \leq Re \leq 2.05 \times 10^4$  and  $4.22 \times 10^3 \leq Re \leq 6.35 \times 10^4$  for the lower and the higher Rayleigh number, respectively. The velocity fields are measured using a 2D2C PIV system, as depicted in Fig. 1b. The PIV System comprises a PCO

SensiCam camera (1280x768pixel, 12bit) and a Continuum Minilite PIV Laser ( $E = 30\text{mJ/pulse}$ ) connected with an ILA Synchronizer. The measurements are carried out at  $f = 2\text{Hz}$ , a measurement time of  $t = 900\text{s}$  and at  $z = 0.5$ . DEHS particle are used for seeding. All measured velocities are normalized to the inlet velocity  $u_{in}$ .

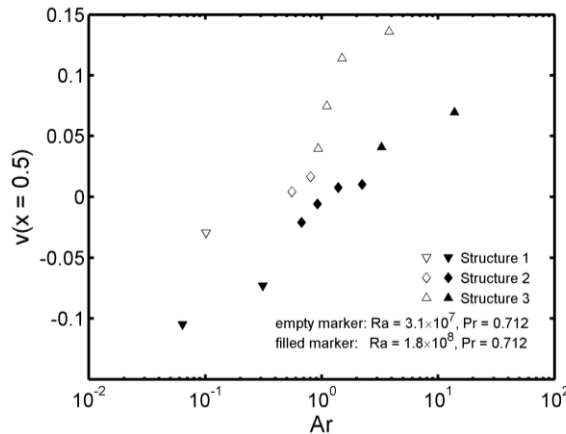
## RESULTS & DISCUSSION

Figure 2 depicts the found flow structures within  $6.4 \times 10^{-2} \leq Ar \leq 1.39 \times 10^1$  using the vertical and horizontal velocity component  $u$  and  $v$  averaged both over time and  $x$  and  $y$ , respectively. The red dots represents the averaged horizontal velocity component  $u(y)$  and the blue dots the averaged vertical component  $v(x)$ . The dash black line shows a simple model on how the structures should look like. In Figure 2a the flow structure is dominated by forced convection, which is why  $v(x \approx 0.5) < 0$  and  $u(x) \approx 0$ . This can be explained by two counter-rotating large eddies as described in [5]. Figure 2b and 2c represents a symmetry break, where one large eddy dominates full domain as described in [4], whereas  $v(x=0.5) \approx 0$ . In Figure 2d the thermal convection is dominant, i.e. the mean flow points upwards. In a two dimensional case the flow structure would be the opposite of 2a. But in the present case one large eddy evolves, which rotates around  $x$ . However, at this case  $v(x=0.5) > 0$ . So there are 3 different structures within the investigated  $Ar$  range.



**Fig. 2** Three different flow structures are found within  $6.4 \times 10^{-2} \leq Ar \leq 1.39 \times 10^1$ ,  $4.22 \times 10^3 \leq Re \leq 6.35 \times 10^4$  and  $Ra = 3.1 \times 10^7$  and  $Ra = 1.8 \times 10^8$ ; **a** structure 1; **b,c** structure 2; **d** structure 3 (blue dots:  $v(x)$ ; red dots:  $u(y)$ ; black dashed line: simple model of flow structures)

Figure 3 shows  $v(x=0.5)$  within  $6.4 \times 10^{-2} \leq Ar \leq 1.39 \times 10^1$  and  $Ra = 3.1 \times 10^7$  (empty symbols) and  $Ra = 1.8 \times 10^8$  (filled symbols). The downward pointing triangles represent the structure 1 (Fig. 2a). The upward pointing triangles represent the structure 3 (Fig. 2d) and the diamonds represent the structure 2 (Fig. 2b,c). The diagram shows a distinct Archimedes number dependency. But, which is more important, there is also a strong dependency regarding the Rayleigh number. Here, the values of  $v(x=0.5)$  at  $Ra = 1.8 \times 10^8$  are lower than the ones for  $Ra = 3.1 \times 10^7$ . A reason for this might be the increased turbulence. Furthermore, the transitions between the structures are shifted to higher  $Ar$  values. So, the flow structures might become more stable with increasing  $Ra$ . However, the main conclusion from these findings is that the Archimedes number  $Ar$  is not the main parameter for the formation of the large-scale flow structures.



**Fig. 3** Archimedes number dependency of the formation of the large-scale flow structure for Rayleigh numbers  $Ra = 3.1 \times 10^7$  and  $Ra = 1.8 \times 10^8$  (structure 1 = Fig. 2a; structure 2 = Fig. 2b,c; structure 3 = Fig. 2d)

## References

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